

The Giant Monopole Resonance in the $^{112-124}\text{Sn}$ Isotopes and the Symmetry Energy Term in Nuclear Incompressibility

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Abstract

We have investigated the isoscalar giant monopole resonance (GMR) in the Sn isotopes, using inelastic scattering of 400-MeV α -particles at extremely forward angles, including 0° . A value of -550 ± 100 MeV has been obtained for the asymmetry term, K_τ , in the nuclear incompressibility.

Key words:

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Incompressibility of nuclear matter remains a focus of experimental and theoretical investigations because of its fundamental importance in defining the equation of state (EOS) for nuclear matter. The compressional-mode giant resonances – the Giant Monopole Resonance (GMR) and the isoscalar giant dipole resonance (ISGDR), an exotic compressional mode of nuclear oscillation – provide a direct means to experimentally determine the nuclear incompressibility. From recent measurements on the GMR and the

ISGDR, a value of $K_\infty = 240 \pm 10$ MeV has been obtained, consistent with results of recent theoretical results in both relativistic and non-relativistic frameworks [1,2,3].

The asymmetry term, K_τ , in the expression for nuclear incompressibility is important because it is critical to determining the radii of neutron stars, and also in understanding the compressional-mode resonances in very neutron-rich nuclei, the primary focus of investigations at current and forthcoming radioactive ion beam facilities. In particular, it has been shown that the radius of a 1–1.5 M_\odot neutron star is mostly determined by the density dependence of the asymmetry-energy term [4,5].

In this presentation, we report on new measurements on GMR in the even-A Sn isotopes. The experiment was performed at the ring cyclotron facility of the Research Center for Nuclear Physics (RCNP), Osaka University, using inelastic scattering of 400-MeV α particles at extremely forward angles, including 0° . Details of the experimental techniques and the data analysis procedures have been provided previously [6]. Sample background-free “ 0° ” inelastic spectra are presented in Fig. 1 (a). Giant monopole resonance strength distributions were obtained for all Sn isotopes under study using a multipole-decomposition analysis (MDA) [6]; Fig. 1 (b) shows a sampling of MDA fits to angular distribution data. Table I lists the centroid energies and widths extracted from Lorentzian-peak fits to the extracted GMR strength distributions, as well as the “standard” moment-ratios typically used in theoretical calculations.

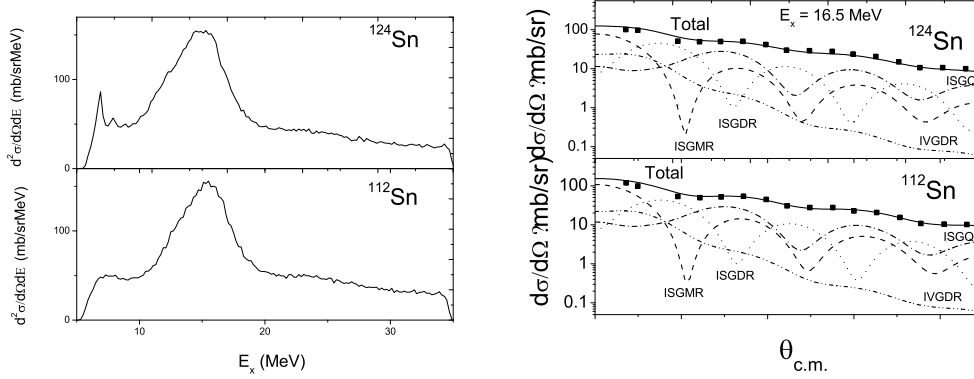


Fig. 1. (a) 0° inelastic scattering spectrum for $^{112}\text{Sn}(\alpha, \alpha')$ and $^{124}\text{Sn}(\alpha, \alpha')$ at $E_\alpha = 400$ MeV. (b) Angular distribution of 1-MeV bins centered at $E_x = 16.5$ MeV for $^{112}\text{Sn}(\alpha, \alpha')$ and $^{124}\text{Sn}(\alpha, \alpha')$. The solid squares are the experimental data and the solid line is the MDA fit to the data. Also shown are the contributions to the fit from $L=0$ (dashed line), $L=1$ (dotted line), $L=2$ (dash-dotted line) and IVGDR (dash-dot-dotted line).

Table 1

Lorentzian-fit parameters and moment-ratios for the GMR strength distributions in the Sn isotopes.

| Target | E_{GMR} (MeV) | Γ (MeV) | m_1/m_0 (MeV) | $\sqrt{m_3/m_1}$ (MeV) | $\sqrt{m_1/m_{-1}}$ (MeV) |
|-------------------|------------------------|----------------|-----------------|------------------------|---------------------------|
| ^{112}Sn | 16.1 ± 0.1 | 4.0 ± 0.4 | 16.2 ± 0.1 | 16.7 ± 0.2 | 16.1 ± 0.1 |
| ^{114}Sn | 15.9 ± 0.1 | 4.1 ± 0.4 | 16.1 ± 0.1 | 16.5 ± 0.2 | 15.9 ± 0.1 |
| ^{116}Sn | 15.8 ± 0.1 | 4.1 ± 0.3 | 15.8 ± 0.1 | 16.3 ± 0.2 | 15.7 ± 0.1 |
| ^{118}Sn | 15.6 ± 0.1 | 4.3 ± 0.4 | 15.8 ± 0.1 | 16.3 ± 0.1 | 15.6 ± 0.1 |
| ^{120}Sn | 15.4 ± 0.2 | 4.9 ± 0.5 | 15.7 ± 0.1 | 16.2 ± 0.2 | 15.5 ± 0.1 |
| ^{122}Sn | 15.0 ± 0.2 | 4.4 ± 0.4 | 15.4 ± 0.1 | 15.9 ± 0.2 | 15.2 ± 0.1 |
| ^{124}Sn | 14.8 ± 0.2 | 4.5 ± 0.5 | 15.3 ± 0.1 | 15.8 ± 0.1 | 15.1 ± 0.1 |

The incompressibility of a nucleus, K_A , may be expressed as:

$$K_A \sim K_{vol}(1 + cA^{-1/3}) + K_\tau((N - Z)/A)^2 + K_{Coul}Z^2A^{-4/3} \quad (1)$$

For a series of isotopes, the difference $K_A - K_{Coul}Z^2A^{-4/3}$ may, thence, be approximated to have a quadratic relationship with the asymmetry parameter, $[(N-Z)/A]$. Such a quadratic fit, using the customary moment ratio $\sqrt{m_1/m_{-1}}$ for the energy of the GMR in calculating the K_A , gives $K_\tau = -550 \pm 100$ MeV; the quoted uncertainty includes the uncertainties in the values of K_A and K_{Coul} , as well as the statistical uncertainties from the fitting procedure.

From the data on the compressional-mode giant resonances, we now have “experimental” values of both K_∞ and K_τ which, together, can provide a means of selecting the most appropriate of the interactions used in EOS calculations. For example, this combination of values for K_∞ and K_τ essentially rules out a vast majority of the Skyrme-type interactions currently in use in nuclear structure calculations [7]. This is borne out by Fig. 2 in which are plotted the values of K_∞ and K_τ for a number of interactions used in both relativistic and non-relativistic calculations. It is clear that nearly all of them fall outside of the “acceptable” region defined by the values obtained in our measurements, leaving a challenge for the theorists to construct appropriate interactions that meet this criterion.

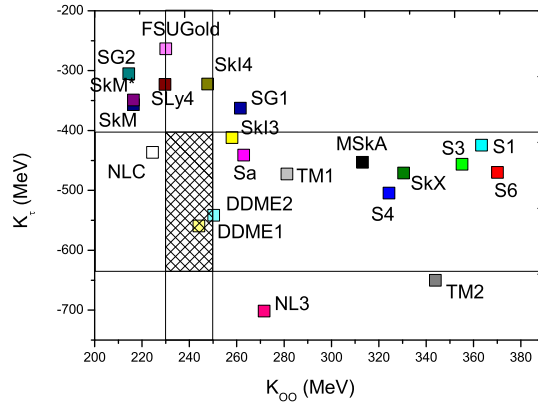


Fig. 2. Values of K_∞ and K_τ calculated from the parameter sets of various interactions as labeled [7]. The vertical and horizontal lines indicate the experimental ranges of K_∞ and K_τ , as determined from the GMR work.

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- References
- [1] G. Colo *et al.*, Phys. Rev. C **70**, 024307 (2004).
 - [2] B. G. Todd-Rutel and J. Piekarewicz, Phys. Rev. Lett. **95**, 122501 (2005).
 - [3] B. K. Agrawal, S. Shlomo and V. Kim Au, Phys. Rev. C **68**, 031304 (2003).
 - [4] J.M. Lattimer and M. Prakash, ApJ **550**, 426 (2001).
 - [5] C.J. Horowitz and J. Piekarewicz, Phys. Rev. Lett. **86**, 5647 (2001).

- [6] M. Uchida *et al.*, Phys. Rev. C **69**, 051301 (2004).
- [7] H. Sagawa *et al.*, arXiv:0706.0966.